Evaluating a Practical Parabolic Time Step Limit for Unconditionally Stable Schemes in a Thermodynamic MHD Model

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ASTRONUM 2023

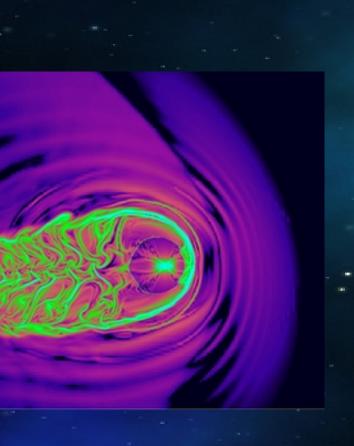
15TH INTERNATIONAL CONFERENCE ON NUMERICAL MODELING OF SPACE PLASMA FLOWS







JUNE 26 – 30, 2023, PASADENA, CA, USA



Unconditionally stable schemes The problem of large time steps Practical time-step limit MAS thermodynamic MHD model Test cases **W** Validation Performance & scaling Summary & future development



- Thermodynamic MHD models (like many others) have multiple time scales leading to vastly different explicit time-step stability requirements
- In order to make simulations *tractable*, we need to exceed the most restrictive limits - here, we focus on the parabolic operators
- Unconditionally stable time stepping schemes are guaranteed to be stable for any sized time step.
- Implicit methods (using iterative matrix solvers)
- Explicit methods (e.g. extended stability Runge Kutta)
- When exceeding explicit time-step limits, one must be careful about accuracy









Exceeding time scales!

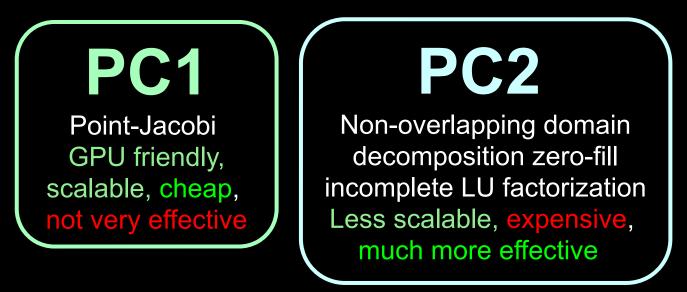
Unconditionally Stable Schemes

Implicit Backward-Euler (BE) + PCG

Backward Euler Ψ

$$\frac{u^{n+1} - u^n}{\Delta t} = F(u^{n+1})$$

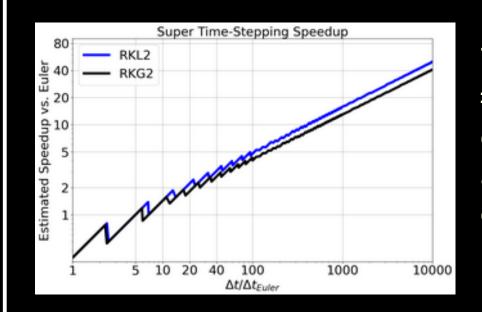
- Yields a system of equations to solve Ψ
- To avoid requiring nonlinear solvers, Ψ we linearize nonlinear terms (e.g. lagged diffusivity)
- We use two non-communicating Ψ preconditioners:



Explicit RKG2(3/2) Super Time-stepping

- Extended Stability Runge-Kutta schemes: Ψ Explicit RK method with stages added for more stability, rather than for more accuracy
- We use the 2nd-order Gegenbauer method Ψ (RKG2) [O'Sullivan (2019)] with an alpha of 3/2

"RKG allows for accurate modeling of solutions at early times and as such if early times are under investigation RKG is the optimal scheme. However, RKC, RKL, RKU, and RKG are all linearly stable and as such will all approach the correct solution asymptotically at long times." [Skaras et. al. (2021)]



TIP: Need to use odd # of stages in RKG2, otherwise amplification factor goes to 1 at highest mode!

Unconditionally Stable Schemes

Implicit BE + PCG

- Robust, proven method
 Can be very efficient
 L-stable
- Can be difficult to implement
- Requires good preconditioner to be efficient, (can be difficult to formulate and implement efficiently)
- * Requires linear(ized) operator
- Solution Solution Solution Solution Solution Solution Content Solution Soluti Sol
- Solution State Accurate Accurate Accurate

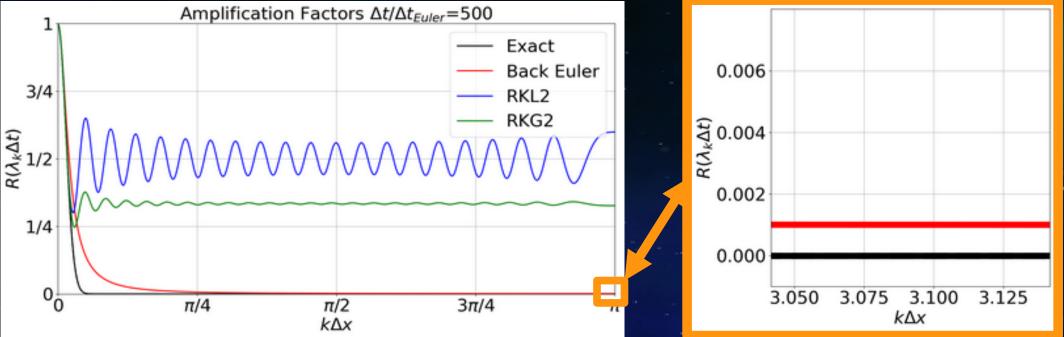


Explicit RKG2(3/2)

- Easy to implement
- ✓ Can include nonlinearities
- ✓ Vectorizable (GPU-friendly)
- ✓ No global synchronization points (better scaling)
- ✓ 2nd-order accurate
- Not as widely adopted & tested
- * Can be slower than implicit methods
- × Only A-stable



 A-stable method can have problems damping high wave modes over limited time scales



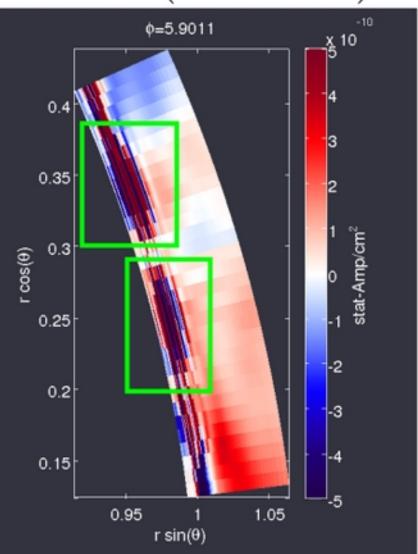
s ly) n points (better scaling)

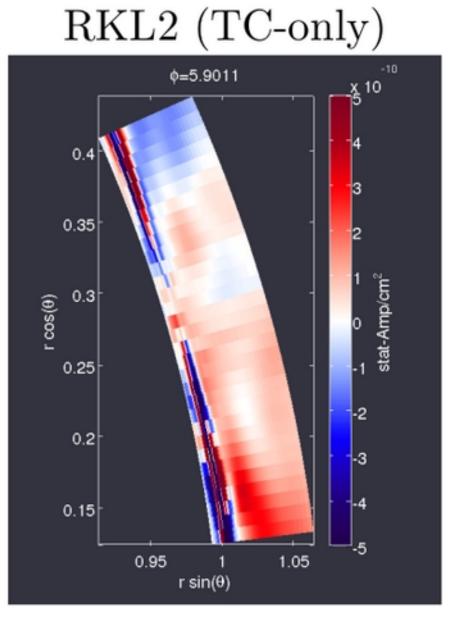
tested cit methods

ASTRONUM 16 Result

4=2.404



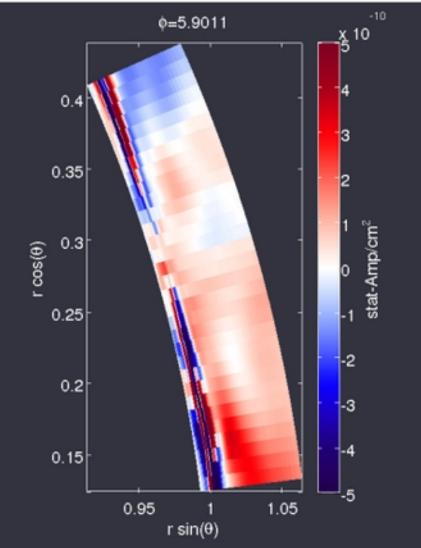




BE+PCG (PC2)

4=2.404

\$=2.404







[Caplan et. al. (2017)]

4=2.404

RKL2 (TC+Visc)

4 -4 -17 -4 0 1 2 3 4 r sin(8)

The Problem with Large Time Steps

Errors with very large time steps are a problem in L-stable methods, but worse in A-stable methods, as they often don't damp high wave modes efficiently

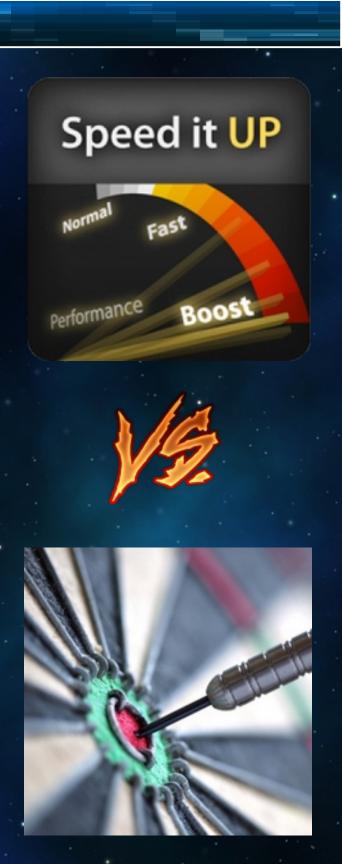
Extended stability Runge-Kutta schemes fall into this latter category, and this issue can limit their applicability and robustness

One option is to run the parabolic advance in a series of "outer" cycles (essentially reducing the time step for the operator) which lowers the errors and repeatedly damps high wave modes

"For general applications the universal approach is to try different numbers of steps and study any sensitivities." [Dawes (2021)]



/ How does adding these cycles affect performance?



A Practical Time Step Limit

Operator split parabolic advance

Discrete form (1st expansion term)

Max abs change in u at grid cell k

Bound the relative change in u at the location of maximum absolute change

$$\frac{\partial \mathbf{u}}{\partial t} = \mathbf{F}(\mathbf{u})$$
$$\frac{u^{n+1} - u^n}{\Delta t} = F(u)$$
$$\left| u^{n+1} - u^n \right| = \Delta t |F(u)|$$

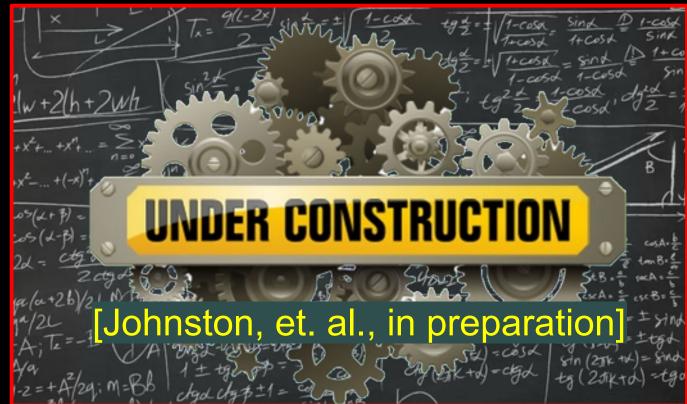
$$F_k(u) \equiv \max\left(|F(u_i)|\right)$$

$$\frac{\left|u_k^{n+1} - u_k^n\right|}{\left|u_k\right|} < \alpha < 1$$

Practical time step limit:
$$\Delta t_p = \alpha \frac{|u_k|}{|F(u_k)|}$$

Applied adaptively: After each outer-cycle, recalculate!

Thermal Conduction: $\mathbf{F}_{\rm tc}(T) = \frac{(\gamma - 1) m_p}{2 k} \frac{1}{\rho} \nabla \cdot (\kappa(T_0) \cdot \nabla T)$ $\kappa(T) = f_{\rm C}(r) f_{\rm mod}(T) \kappa_0 T^{5/2} \, \hat{\mathbf{b}} \hat{\mathbf{b}}$ $f_{\text{mod}}(T) = \left(1 + (T/T_{\text{cut}})^{-10}\right)^{1/4}$



Viscosity: $\mathbf{F}_{\text{visc}}(\mathbf{v}) = \frac{1}{\rho} \nabla \cdot (\nu(\mathbf{x}) \rho \nabla \mathbf{v})$ Lagged diffusivity

Purpose:

General-purpose simulations of the corona and heliosphere for use with solar physics and space weather research

Model:

Spherical resistive thermodynamic MHD

Algorithms:

VEBINAR Implicit & explicit timestepping with finite-difference stencils. Implicit steps use a sparse matrix preconditioned iterative solver

Code:

~70,000 lines of Fortran parallelized with MPI + OpenACC + StdPar

Accelerating a Production Solar MHD Code with Fortran Standard Parallelism

Ronald M. Caplan Predictive Science Inc.

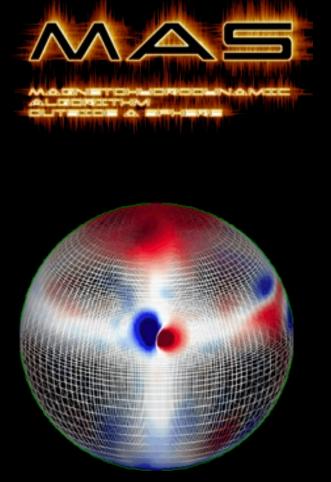
The MAS Code

predsci.com/mas

OpenACC

WEBINAR DIGITAL EVENT JULY 11, 2023 1 PM EDT/10 AM PDT

The MAS MHD Model



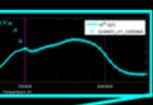
$$\begin{split} \mathbf{A}_{tt} &= \mathbf{v} \times (\nabla \times \mathbf{A}) - \begin{bmatrix} \frac{c^2 \eta}{4\pi} \nabla \times \nabla \times \mathbf{A} \\ \frac{\partial \rho}{\partial t} &= -\nabla \cdot (\rho \mathbf{v}) \\ \hline \mathbf{V}_{tt}^T &= -\nabla \cdot (T \mathbf{v}) - \mathbf{v}^{-2} (T \nabla \cdot \mathbf{v}) + \frac{\gamma - 1}{2k} \frac{m_p}{\rho} \begin{bmatrix} -\nabla \cdot (\mathbf{q}_1 + \mathbf{q}_2) \\ \nabla \cdot (\mathbf{q}_1 + \mathbf{q}_2) \\ \mathbf{v}_{tt}^2 &= -\nabla \cdot (T \mathbf{v}) \\ \hline \mathbf{U}_{tt} &= -\nabla \cdot (\mathbf{v} \cdot \mathbf{v}) + \frac{\gamma - 1}{2k} \frac{m_p}{\rho} \begin{bmatrix} -\nabla \cdot (\mathbf{q}_1 + \mathbf{q}_2) \\ \nabla \cdot \mathbf{v}_{tt} \\ \mathbf{q}_2 &= (1 - f(r)) \frac{k}{(\gamma - 1)} \frac{\rho}{m_p} T \mathbf{v} \hat{\mathbf{b}} \hat{\mathbf{b}} \\ \hline \mathbf{v}_{tt}^2 &= -\nabla \cdot (\epsilon_{\pm} [\mathbf{v} \pm \mathbf{v}_{\mathbf{A}}]) - \frac{c_{\pm}}{2} \\ \hline \mathbf{v}_{tt} &= -\mathbf{v} \cdot \nabla \mathbf{v} + \frac{1}{\rho} \begin{bmatrix} 1 \\ c \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \left(\frac{\epsilon_{+} + \epsilon_{-}}{2} \right) \\ -\nabla \cdot \mathbf{v} &= -\mathbf{v} \cdot \nabla \mathbf{v} + \frac{1}{\rho} \begin{bmatrix} 1 \\ c \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \left(\frac{\epsilon_{+} + \epsilon_{-}}{2} \right) \\ -\nabla \cdot \mathbf{v} &= -\mathbf{v} \cdot \nabla \mathbf{v} + \frac{1}{\rho} \begin{bmatrix} 1 \\ c \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \left(\frac{\epsilon_{+} + \epsilon_{-}}{2} \right) \\ -\nabla \cdot \mathbf{v} &= -\frac{2kT\rho/m_p}{2kT\rho/m_p} \begin{bmatrix} \frac{\partial z_{\pm}}{\partial t} \\ \frac{\partial z_{\pm}}{\partial t} \\ + \frac{z_{\pm}}{4} (\mathbf{v} \mp \mathbf{v}_{\mathbf{A}}) \cdot \nabla (\ln \rho) + \frac{z_{\mp}}{2} (\mathbf{v} \mp \mathbf{v}_{\mathbf{A}}) \cdot \nabla (\ln |\mathbf{v}_{\mathbf{A}}|) \\ \hline \mathbf{v} &= -\frac{2kT\rho/m_p}{2kT\rho/m_p} \begin{bmatrix} \frac{\partial z_{\pm}}{\partial t} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ -\frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ -\frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \hline \mathbf{v} &= -\frac{2kT\rho/m_p}{2kT\rho/m_p} \begin{bmatrix} \frac{\partial z_{\pm}}{z_{\mp}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ -\frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \hline \mathbf{v} &= -\frac{2kT\rho/m_p}{2kT\rho/m_p} \begin{bmatrix} \frac{\partial z_{\pm}}{z_{\mp}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ -\frac{\partial z_{\pm}} \end{bmatrix} \\ \hline \mathbf{v} &= -\frac{2kT\rho/m_p}{2kT\rho/m_p} \begin{bmatrix} \frac{\partial z_{\pm}}{z_{\mp}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}}{z_{\mp} \nabla \times \mathbf{B}} \\ \frac{\partial z_{\pm}} \\ \frac{\partial z_{\pm}}}{z_{\mp} \nabla \times \mathbf{B}} \\ \end{bmatrix}$$

 (r_i, θ_j, ϕ_k)

CORONA

HELIOSPHERE

Different components of the model are used depending on the domain and use-case

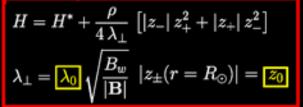


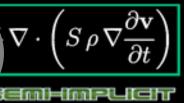


 $(C_{w}^{0}/(1-C_{f})^{2}-1)$

 $\hat{k}^2 (v_c^2 + |v_A|^2)$ + $(r \Delta \theta)^{-2} + (r \Delta \phi \sin \theta)^{-2})$







OPERATOR

$B/ B = B/\sqrt{4\pi\rho} = 6.09 G \gamma p/\rho$	$\begin{split} \beta_{\rm int}(T) &= \left\{ \begin{array}{ll} (T/T_{\rm int})^{-5/2} & T < T_{\rm int} \\ 1, & T \geq T_{\rm int} \end{array} \right. \\ T_{\rm int} &= 3.5 \times 10^6 K \\ f(r) &= 1 - 0.5 \tanh[(r-10R_0)/R_0] \end{split}$	$\begin{split} S &= (\Delta t^2 \bar{k}^2) \\ C_f &= \Delta t \bar{k} \cdot \mathbf{v} \\ C_{\mathbf{w}}^2 &= 0.25 \Delta t^2 \\ \bar{k}^2 &= 4 (\Delta r^{-2} + 1)^2 \end{split}$

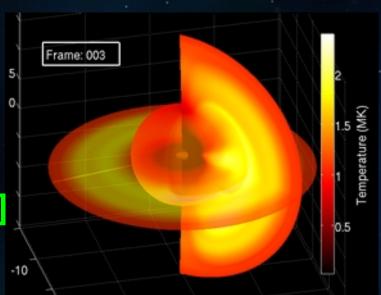


Test Cases

Test 1:

Modified low-res test case based on simulations used in [Reeves et. al. (2019)]

Resolution: 151x151x151 ~ 3.4 million points



Thermodynamic MHD relaxation

Test 2:

Modified test case used in Astronum 2016 paper [Caplan et. al. (2017)]



Resolution: 181x251x502 ~ 22.8 million points

Validation runs:

Integrate relaxation for ~8 simulation-hours

Scaling runs (Test 2 only):

Integrate for 6 simulation-minutes starting with the ~8 simulation-hour relaxation (subtract restart loading time)

Thermodynamic MHD relaxation



Testing Procedure

Questions to address:

□ Can we get a solution as good (or better) as BE+PCG with RKG2 if we automatically outer-cycle at the practical time step limit? Can we get a "better" solution with BE+PCG if we outer-cycle it? How does outer-cycling affect performance? □ Is RKG2 competitive with BE+PCG?

Tests:

SC1: 1 outer cycle (original) **SCA**: Automatic adaptive outer cycles (practical time step)









Computational Environment

Tests performed on both
 CPUs and GPUs

- On CPUs, the PC2 preconditioner is used for BE+PCG; on GPUs, PC1 is used
- On CPUs, the gfortran compiler is used with OpenMPI 4, on GPUs, the NV compiler is used with OpenMPI 3

In-house workstation with NVIDIA RTX 3090 Ti



CPU	Core i5-13600KF
GPU	RTX 3090 Ti
Peak DP FLOPs	0.625 TFLOP/s
Memory	24 GB
Memory Bandwidth	1008 GB/s

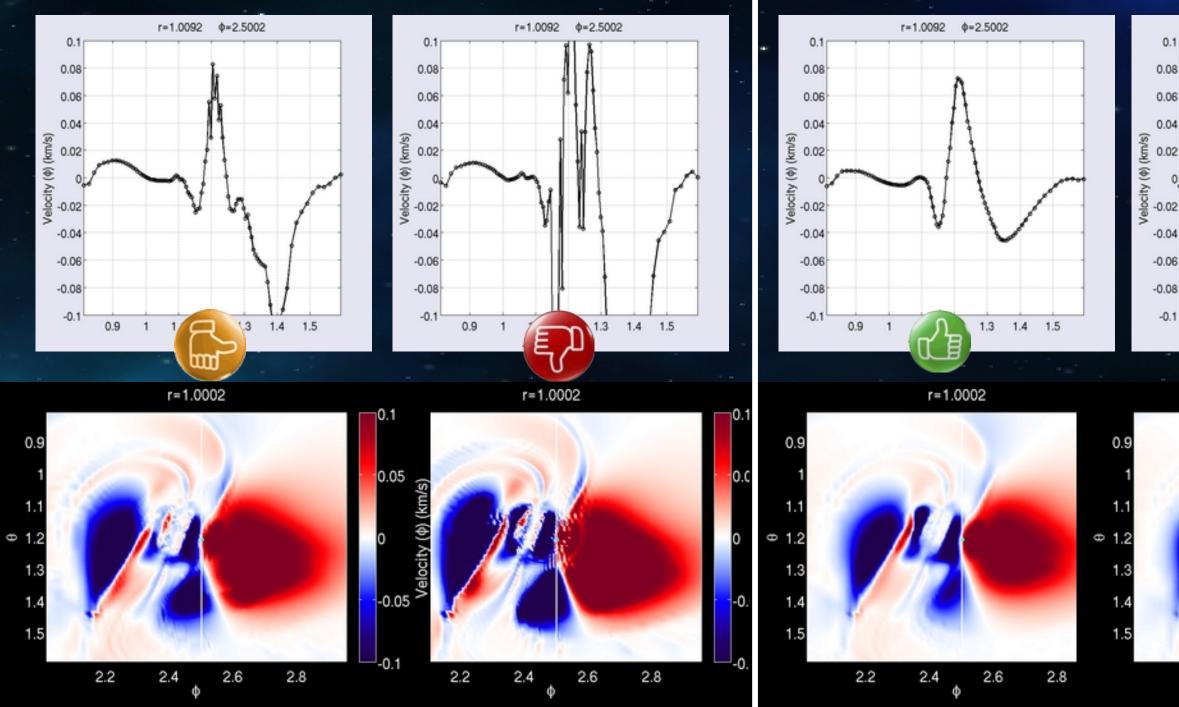


Validation Results Test 1

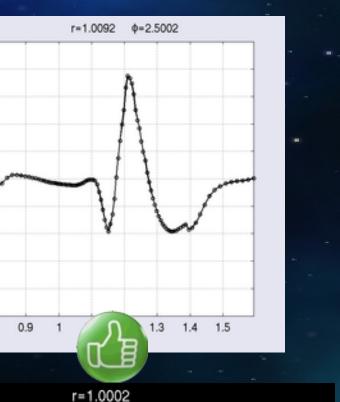
RKG2 SC1

BE+PCG SC1

BE+PCG SCA

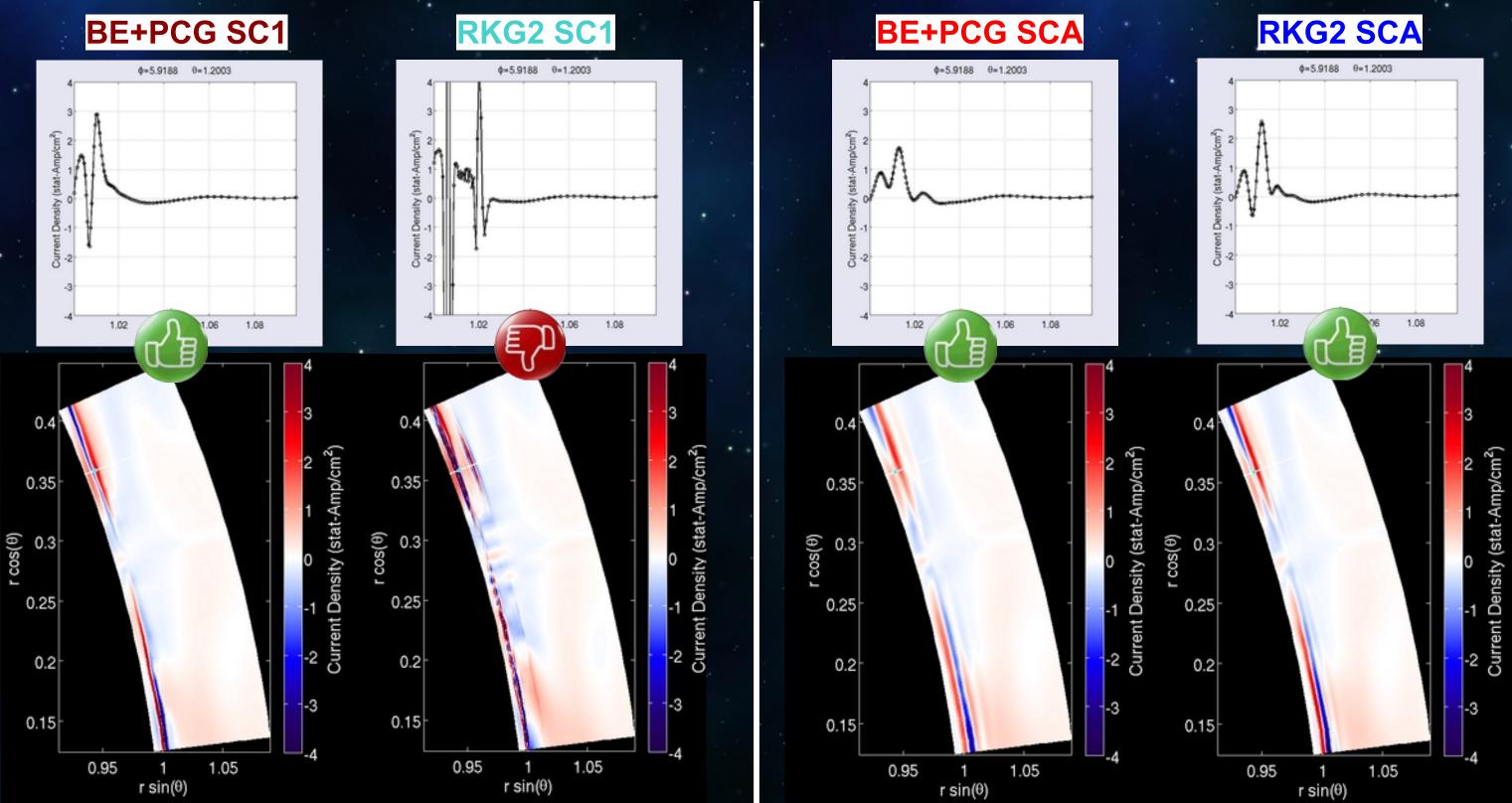


RKG2 SCA

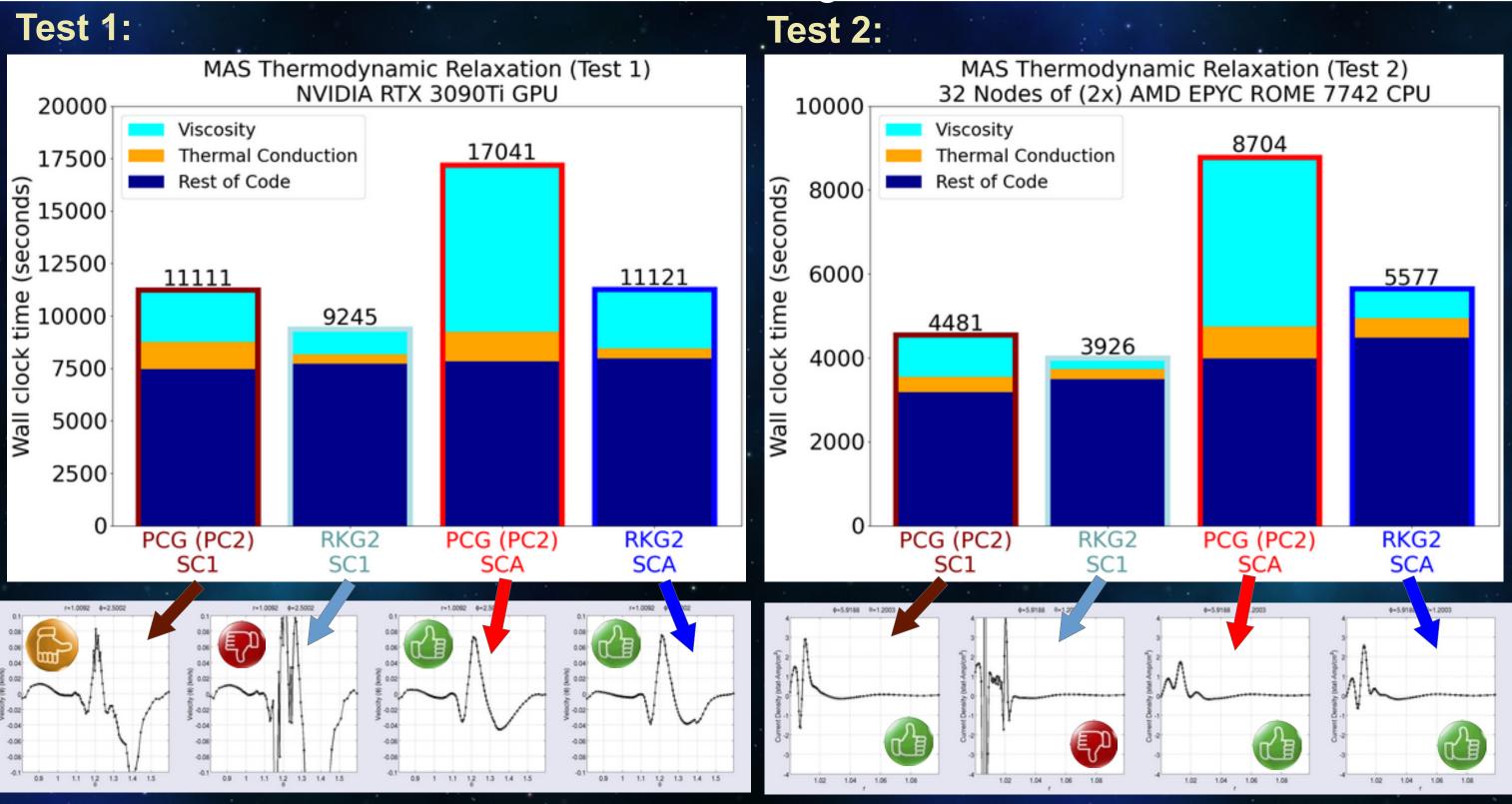


2.2 2.4 2.6 2.8 0.1 0.1 0.05 (s/w) (4) Aijona, -0.1 -0.1

Validation Results Test 2

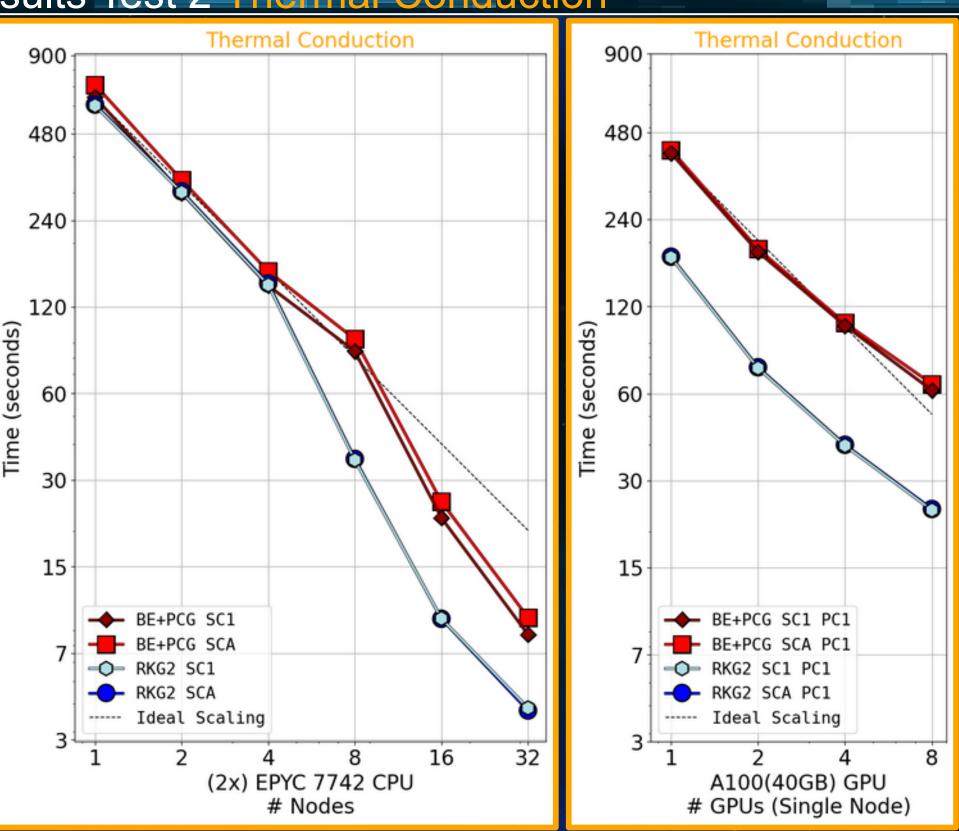


Validation Timing Results



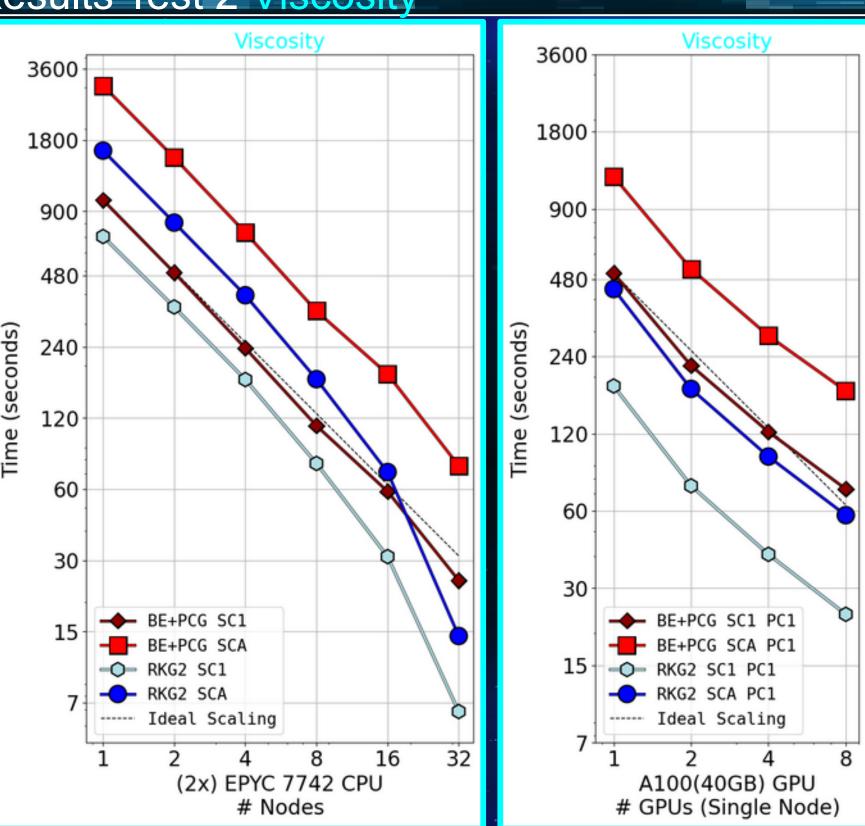
Scaling Results Test 2 Thermal Conduction

- Very few outer cycles were needed in this test for TC
- Therefore, SCA has similar run times as SC1
- Scaling is better with RKG2 than BE+PCG (PC2) on CPUs
- W Run time for RKG2 is much faster than BE+PCG (PC1) on GPUs
- Scaling of both schemes similar on multi-GPU single server runs



Scaling Results Test 2 Viscosity

- The test problem required ~10 outer cycles for viscosity
- Therefore, auto-cycle has slower run time than single cycle
- Scaling is better with RKG2 than
 PCG (PC2) on CPUs
- Run time for RKG2 is faster than PCG (PC1) on GPUs with similar scaling
- RKG2 auto-cycle faster than PCG single cycle:
 - With PCG (PC1), always the case
 - With PCG (PC2), it's faster with maximum CPUs due to better scaling and PC2 effectiveness decline

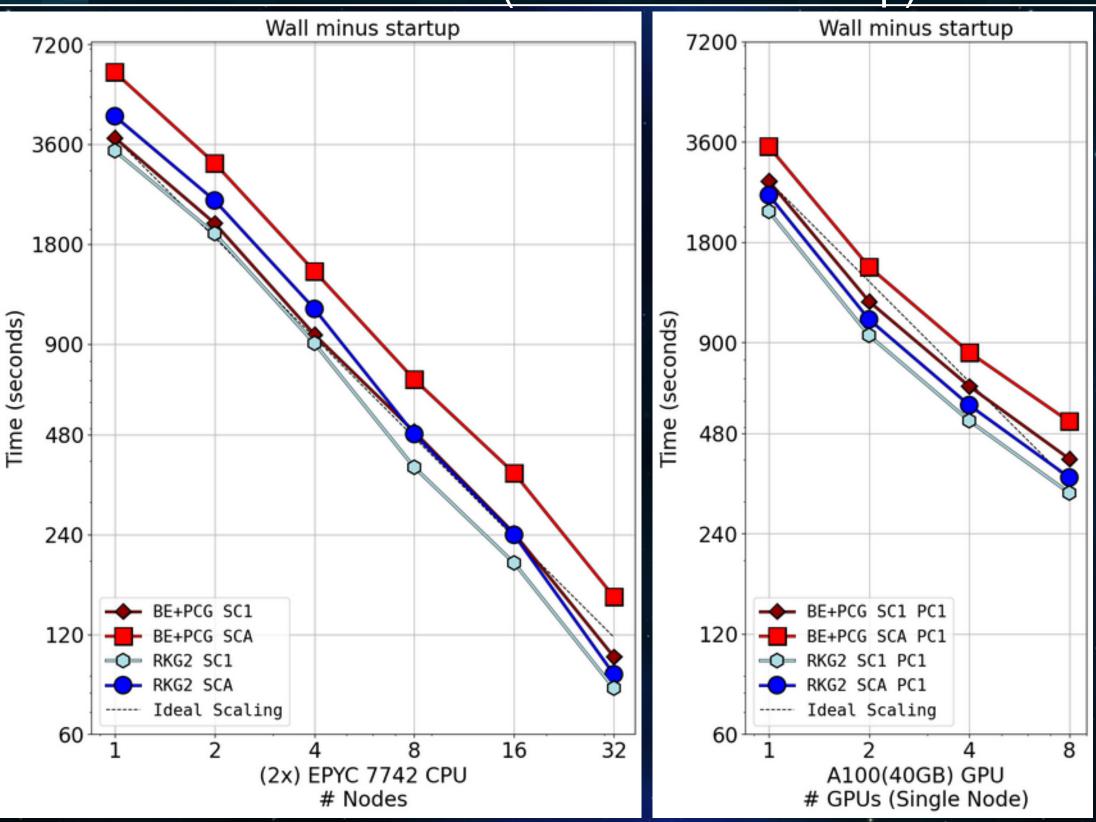


Scaling Results Test 2 Total Wall Clock (minus restart startup)

RKG2 auto-cycle has comparable or better performance and scaling than **BE+PCG** singlecycle

Ψ

BE+PCG auto-cycle Ψ significantly slower than single-cycle





Results

□ Can we get a solution as good (or better) as BE+PCG with RKG2 if we automatically outer-cycle at the practical time step limit? Yes (in our cases) □ Can we get a "better" solution with BE+PCG if we outer-cycle it? Yes (in our cases) How does outer-cycling affect performance? For BE+PCG, significant decrease in performance (in our cases) For RKG2, small decrease in performance (in our cases) □ Is RKG2 competitive with BE+PCG? Yes (in our cases)

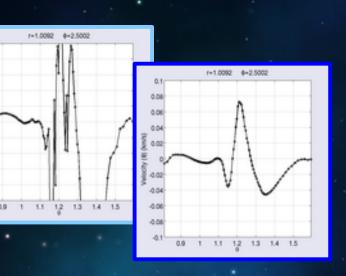
How general are these results???



Summary

- Unconditionally stable methods are often necessary when conditionally stable schemes' time-step limits for operators (e.g. parabolic) are too restrictive for practical simulations
- Using "too large" of a time step can create large errors including qualitative changes in the operator's effect and/or inability to damp high wave modes (oscillations)
- We have tested an easy-to-calculate practical time step limit for such schemes that can be applied as an operator-split outer cycle
- W The new limit can help fix known solution issues in extended stability Runge-Kutta (super time stepping) methods that are due to their poor amplification factors at high wave modes
- Y Testing the method with RKG2 implemented in the MAS MHD code recovered the solution accuracy of the BE+PCG method with negligible effect on performance & scaling
- Work is proceeding on the theoretical formulation/modification of the practical time step, including careful testing, especially in nonlinear cases [Johnston et al, in preparation]

 $\Delta t_p = \alpha$ $k: |F(u_k)| = \max |F(u)|$



MAS Thermodynamic Relaxation (Test 1) NVIDIA RTX 3090Ti GPU

15000

12500 10000

7500

5000

2500

